



Towards Natural Tuning

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Principal Technical Result

NB: Construction will be in (1+1)d. 2d theories are special in many respects, but not as far as the hierarchy problem goes

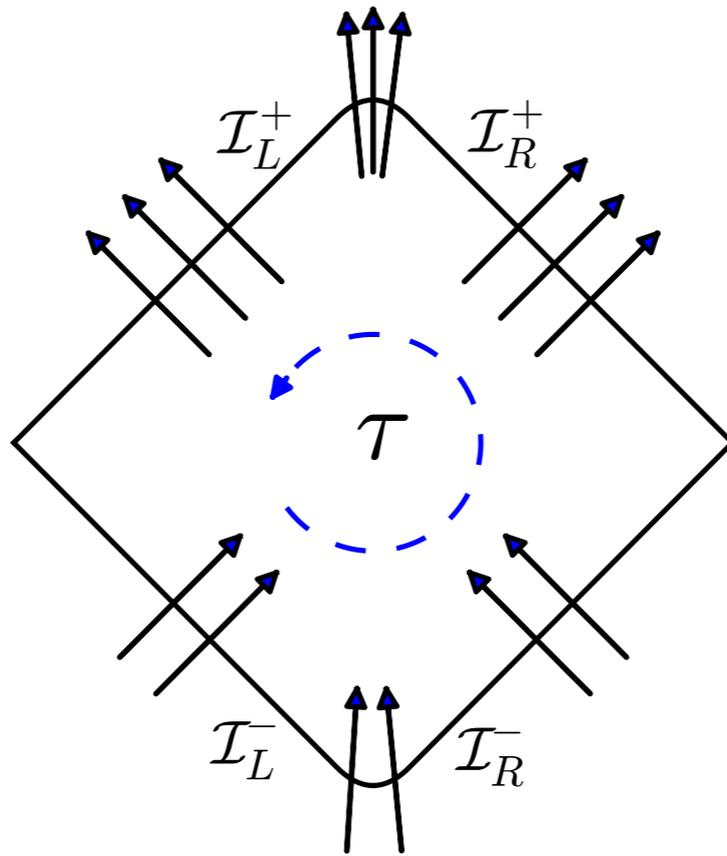
Start with an arbitrary UV complete natural QFT $\mathcal{L}(\psi, H)$
Non-protected scalars are allowed as soon as they are heavy



Calculate S-matrix $S_n(p_i)$



“Gravitational dressing” gives $\hat{S}_n(p_i, \ell)$



$$\hat{S}_n(p_i) = e^{i\ell^2/4 \sum_{i < j} p_i^* p_j} S_n(p_i)$$

Properties of gravitational dressing

- * Results in a well-to-do S-matrix
- * Physical spectrum remains the same
- * Low energy EFT description, tuned for $m\ell \ll 1$

$$\mathcal{L}(\psi, H) + \sum_{\Delta_i > 2} \ell^{\Delta_i - 2} \mathcal{O}_i$$

free massive scalar:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 + \frac{\ell^2}{8} \left((\partial\phi)^4 - m^4\phi^4 \right) + \dots$$

***THIS CONSTRUCTION SHOULD NOT
BE POSSIBLE !!!**

This is how these theories should have been found:

What are possible integrable reflectionless massless theories in two dimensions?

Everything is determined by a two-particle phase shift:

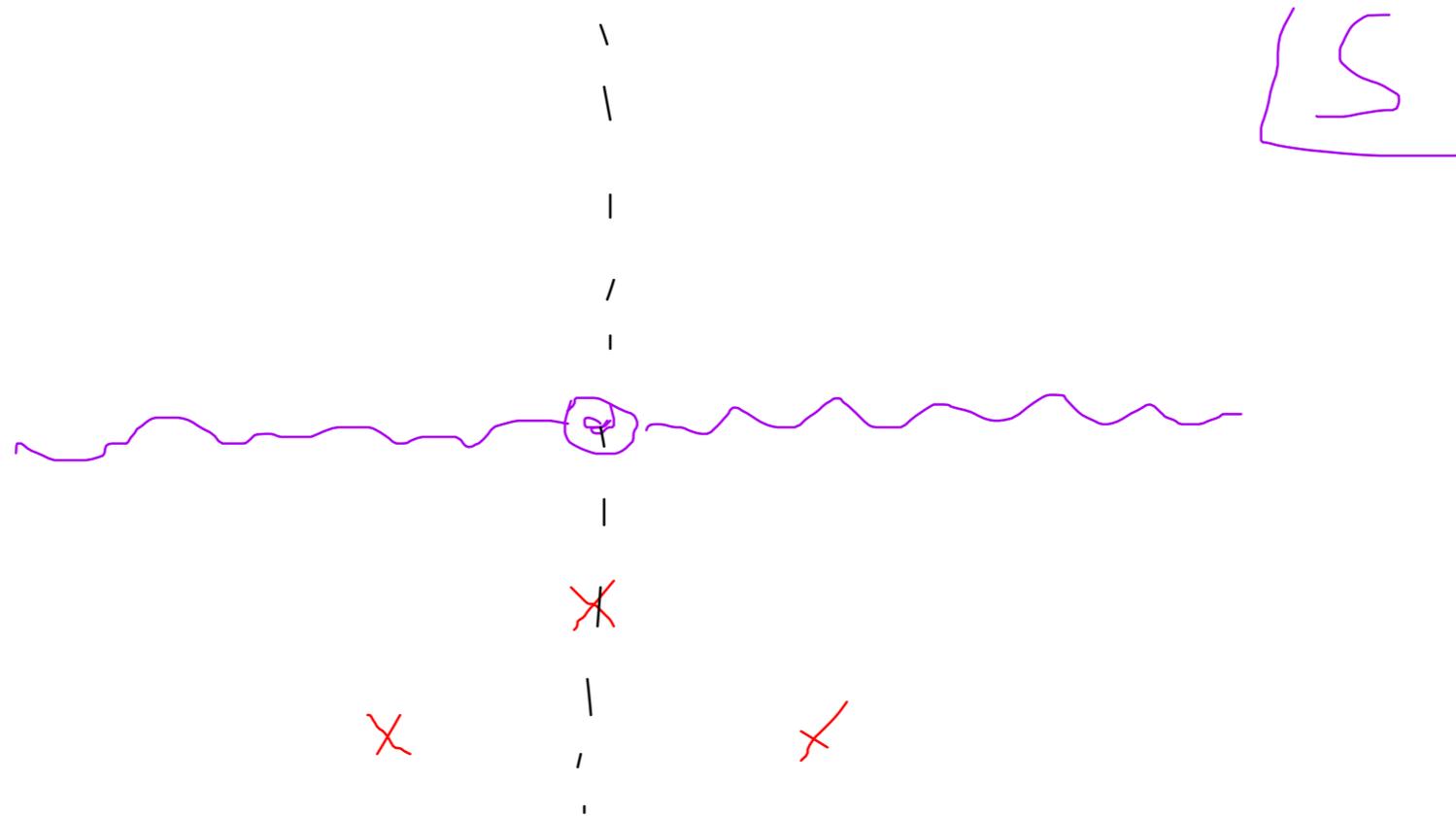
$$S = e^{2i\delta(s)} \mathbf{1}$$

Unitarity+Analyticity+Crossing:

Zamolodchikov '91

$$e^{2i\delta(s)} = \prod_j \frac{\mu_j + s}{\mu_j - s} e^{iP(s)}$$

$$\text{Im } s > 0$$



Expectation from Locality: $P(s) = 0$

Goldstino (Volkov-Akulov) Theory

$$\mathcal{L} = \psi \bar{\partial} \psi + \bar{\psi} \partial \bar{\psi} - \frac{1}{M^2} (\psi \partial \psi) (\bar{\psi} \bar{\partial} \bar{\psi}) + \dots$$

$$e^{2i\delta_{\text{Gold}}(s)} = \frac{iM^2 - s}{iM^2 + s}$$

A simple example of “Asymptotic Safety”:

naively non-renormalizable theory flows into a strongly coupled UV fixed point, no new stuff added

Corresponds to integrable RG flow between tricritical Ising model in the UV and Ising model in the IR

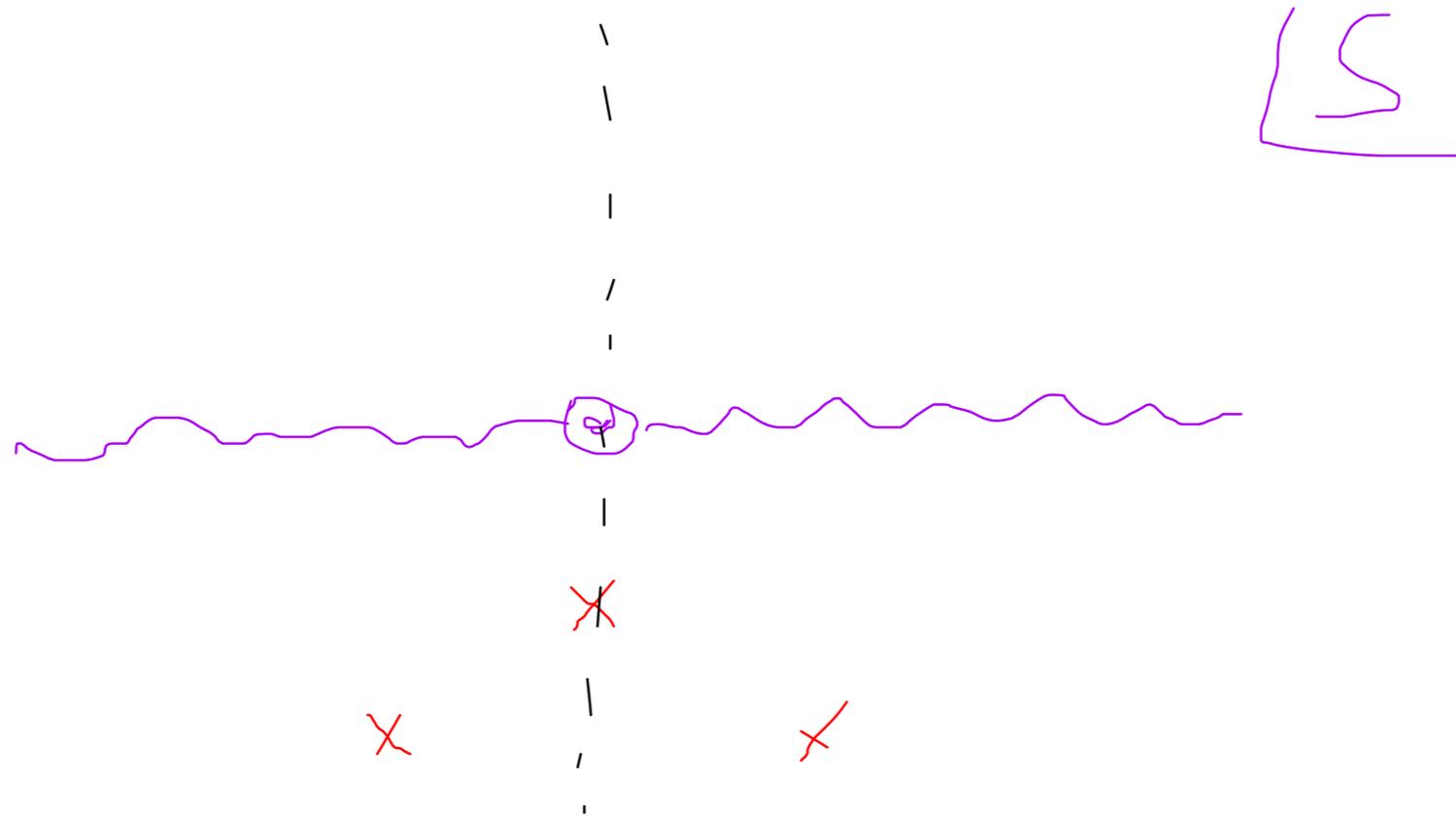
(equivalently, N=1 Wess-Zumino model in the UV and free fermion in the IR)

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$$\text{Im } s > 0$$



Expectation from Locality: $P(s) = 0 + \ell^2 s$

Let us look at
at (D-2) bosons with

$$e^{2i\delta(s)} = e^{is\ell^2/4}$$

- *Polynomially bounded on the **physical** sheet
- *No poles anywhere. A cut all the way to infinity with an infinite number of broad resonances
- *One can reconstruct the entire finite volume spectrum using Thermodynamic Bethe Ansatz

$$E(N, \tilde{N}) = \sqrt{\frac{4\pi^2(N - \tilde{N})^2}{R^2} + \frac{R^2}{\ell^4} + \frac{4\pi}{\ell^2} \left(N + \tilde{N} - \frac{D-2}{12} \right)}$$

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A new type of RG flow behavior:

Asymptotic Fragility

Integrable theory of gravity

* (

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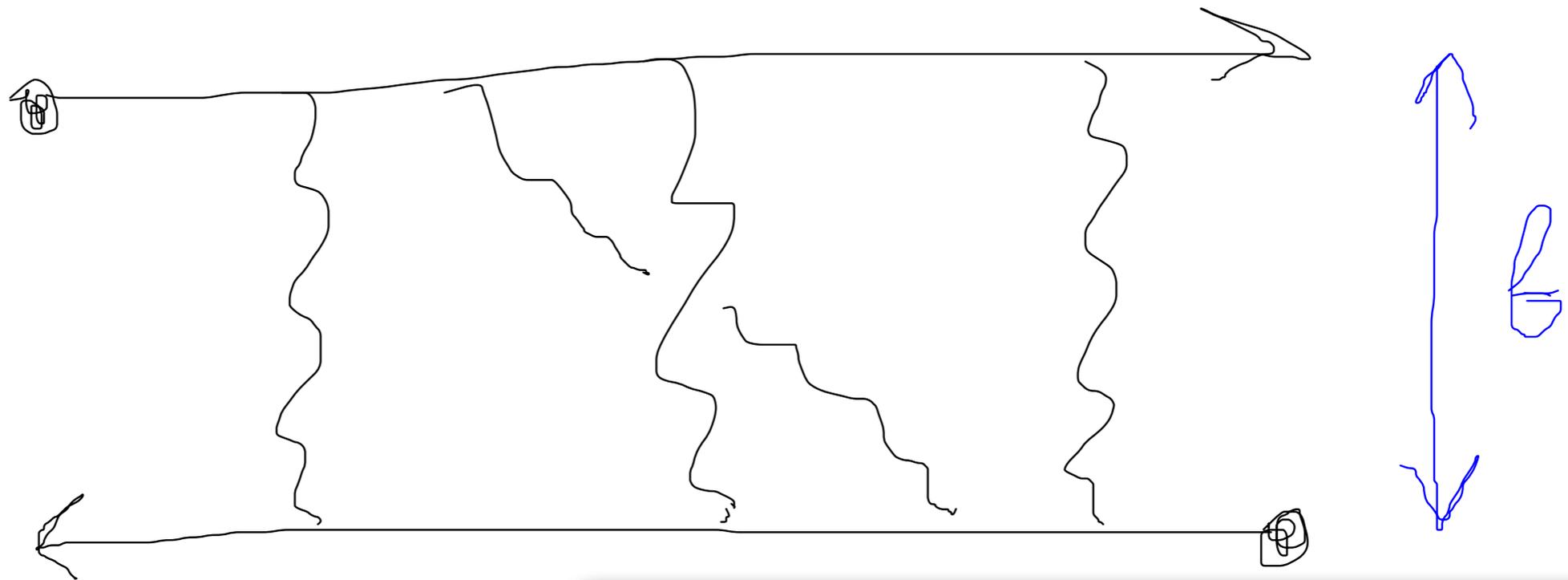
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Integrable QG rather than QFT

Gravitational shock waves:

Dray, 't Hooft '85
Amati, Ciafaloni, Veneziano '88

$$S \gg M_{\text{pl}}^2$$
$$b \gg R_s$$



Eikonal phase shift:

$$e^{i2\delta_{\text{eik}}(s)} = e^{i\ell^2 s/4}$$

$$\ell^2 \propto G_N b^{4-d}$$

Some properties of the theory

classical action:

$$S_{NG} = -\ell^2 \int d^2\sigma \sqrt{-\det(\eta_{\alpha\beta} + \partial_\alpha X^i \partial_\beta X^i)}$$

- * Theory of gravitational shock waves.
- * No UV fixed point and central charge.
- * Maximal achievable (Hagedorn) temperature.
- * Integrable cousins of black holes.
- * Minimal length.
- * No local off-shell observables.

Integrable Black Hole Precursors

Time Delay

$$\Delta t_{cms} = \frac{1}{2} \ell_s^2 E_{cms}$$

c.f. $\Delta t_H = \ell_{Pl}^4 E_{cms}^3$ for Hawking evaporation in 4d

Equivalence Principle at work

Δt is the same for a single hard particle and for a bunch of soft ones

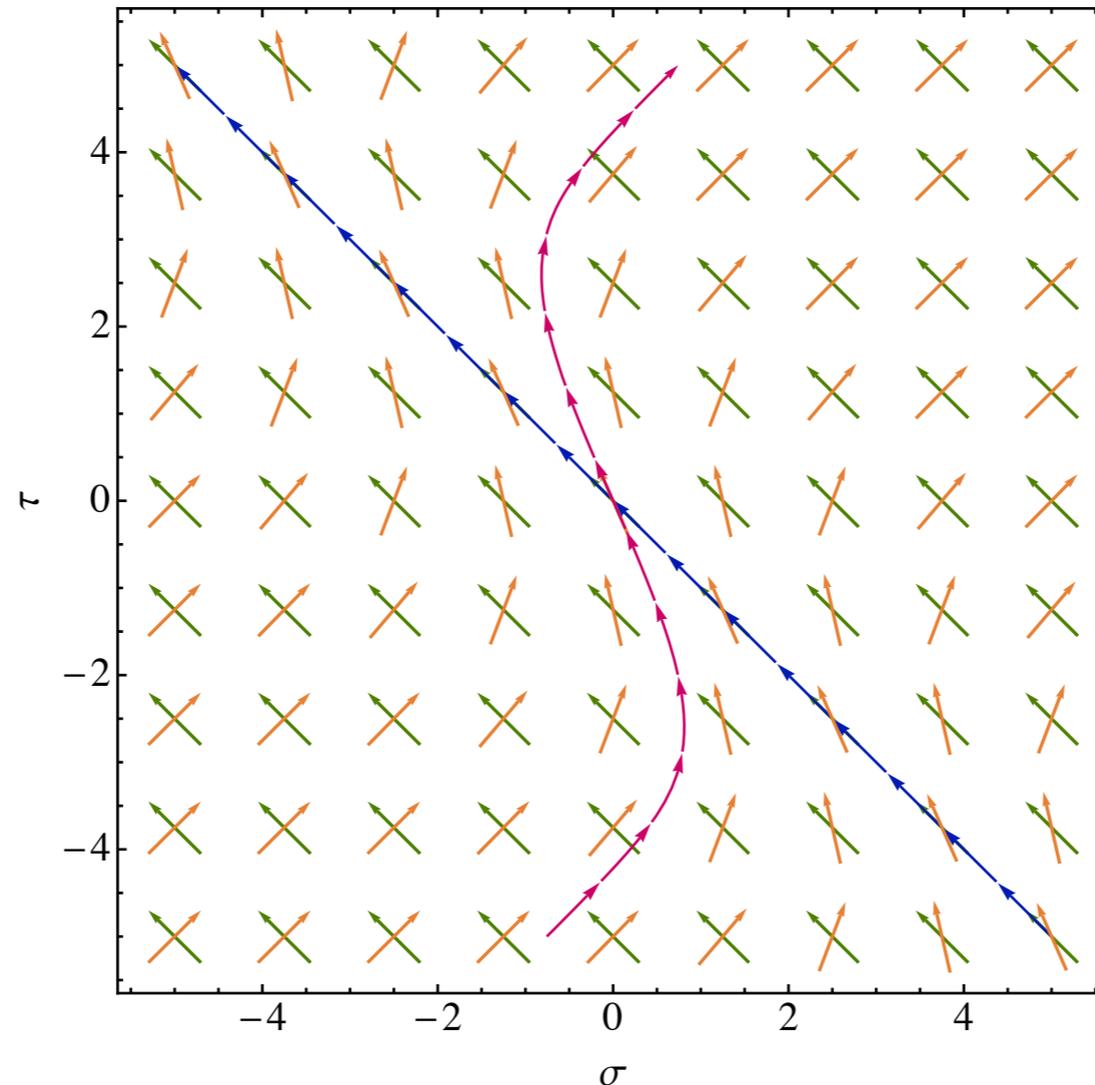
String uncertainty principle

$$\Delta x_L \Delta x_R \geq \ell_s^2$$

for identical packets $\Delta x_{out}^2 = \Delta x_{in}^2 + \frac{\ell_s^4}{\Delta x_{in}^2}$

Classical Origin of the Time Delay

$X_{cl}^i(\tau + \sigma)$ is a solution



$$\Delta t = \int_{-\infty}^{\infty} dz X_{cl}'^2 = \ell^2 E$$

exactly reproduces the quantum answer

*This was an integrable QG coupled to
(D-2) massless bosons.*

*Is there a generalization to other (non-integrable)
theories?*

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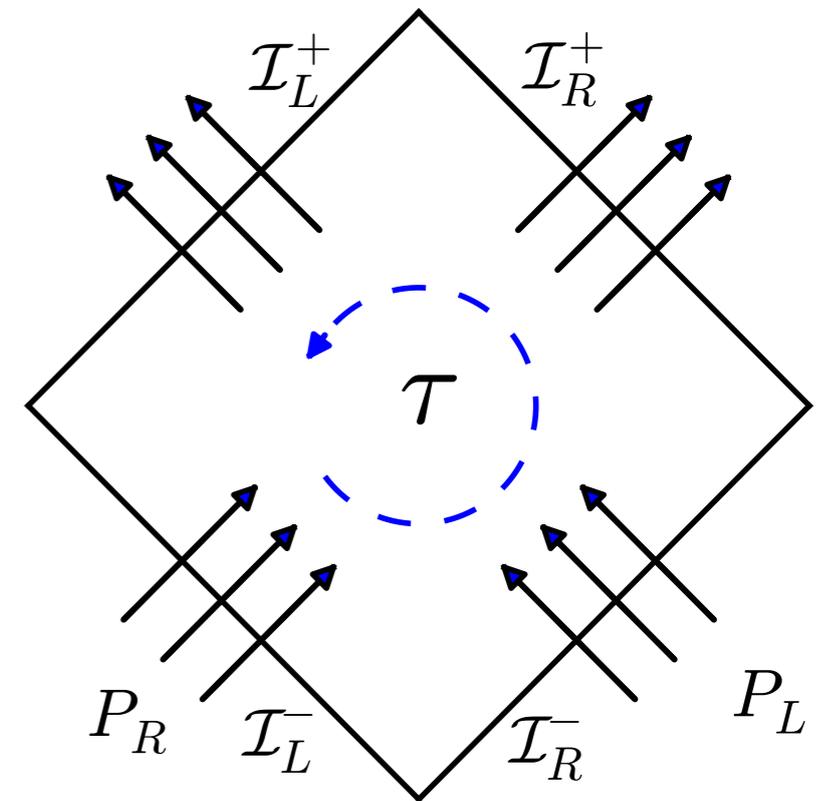
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Eikonal Scattering From Boundary Quantum Mechanics

Verlinde & Verlinde '91

$$g_{\mu\nu} = \begin{pmatrix} e^\phi \partial_\alpha X^a \partial_\beta X^b \eta_{ab} & 0 \\ 0 & h_{ij} \end{pmatrix}$$

$$S_{CS}[X] = \ell^{-2} \oint d\tau \epsilon_{\alpha\beta} X^\alpha \partial_\tau X^\beta$$



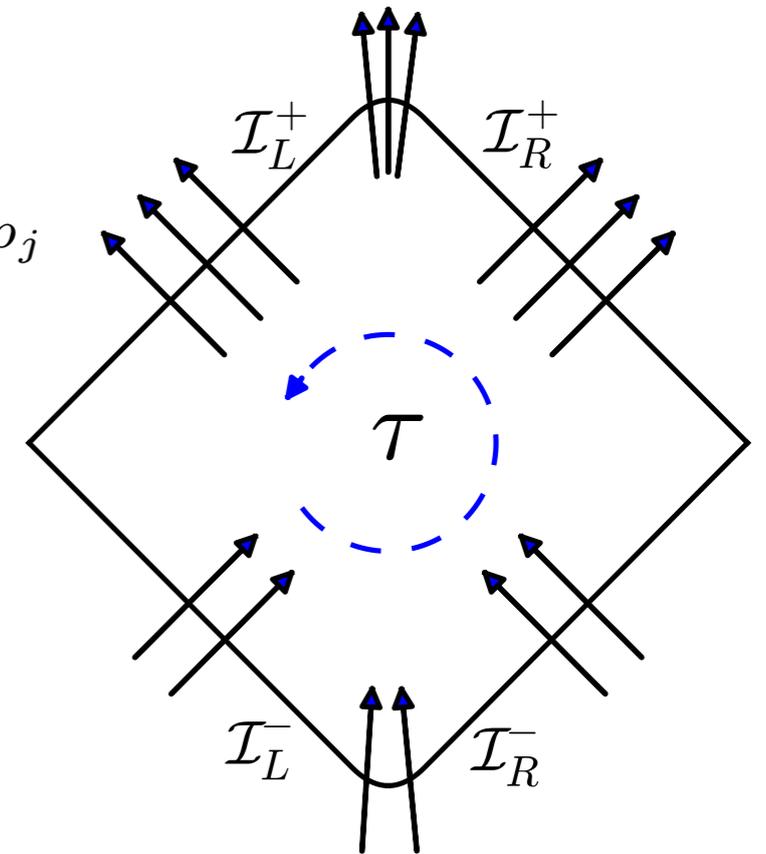
$$S_{eik} = \int \mathcal{D}X e^{iS_{CS}[X] + i(\sum_i p_{iR}^\alpha X_\alpha(\tau_i) + \sum_j p_{jL}^\alpha X_\alpha(\tau_j) + \sum_i \bar{p}_{iR}^\alpha X_\alpha(\bar{\tau}_i) + \sum_j \bar{p}_{jL}^\alpha X_\alpha(\bar{\tau}_j))}$$

Most simple-minded generalization:

$$\mathcal{D}(p_i) = \int \mathcal{D}X e^{iS_{CS}[X] + i \sum_i p_i^\alpha X_\alpha(\tau_i)} = e^{i\ell^2/4 \sum_{i<j} p_i * p_j}$$

$$p_i * p_j = \epsilon_{\alpha\beta} p_i^\alpha p_j^\beta$$

does not produce a consistent S-matrix,
but allows to dress:



$$\hat{S}_n(p_i) = e^{i\ell^2/4 \sum_{i<j} p_i * p_j} S_n(p_i)$$

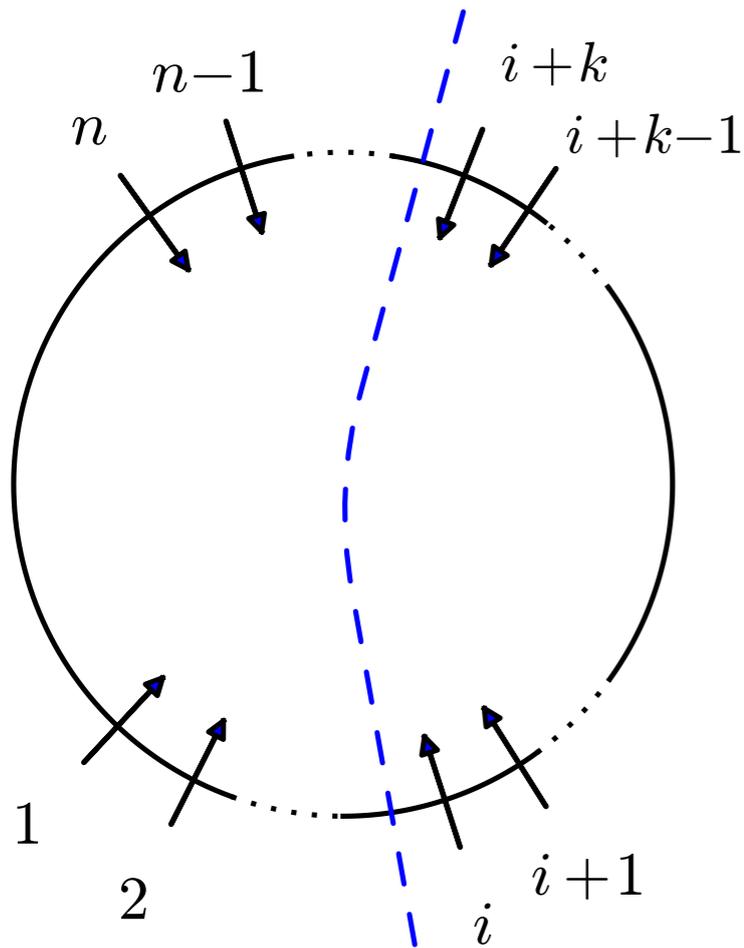
✓ Crossing Symmetry

✓ Analyticity

? Unitarity

? Factorization

✓ Unitarity from Factorization



$$\mathcal{D}(p_i) = e^{i\ell^2/4 \sum_{a < a'} k_a * k_{a'}} e^{i\ell^2/4 \sum_{b < b'} q_b * q_{b'}}$$

$$\hat{S} = U S U$$

$$U |\{k_a\}\rangle = e^{i\ell^2/4 \sum_{a < a'} k_a * k_{a'}} |\{k_a\}\rangle$$

The whole story is a bit similar to non-commutativity.

Two crucial differences:

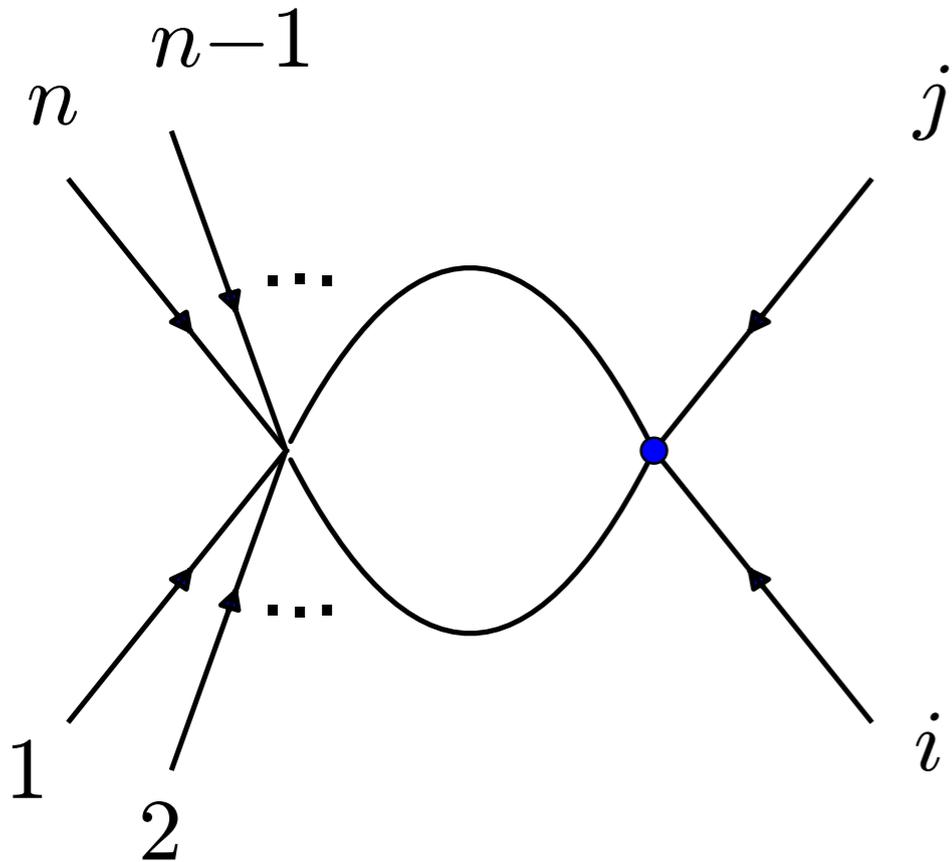
- * Dressing of the full S-matrix, rather than of the tree amplitudes.
- * No summation over different cycling orderings. Preserves causality.

Perturbative Check

$$\mathcal{L}_{QFT} = \frac{1}{2} \left(\sum_i \partial\phi_i \partial\phi_i - m_i^2 \phi_i^2 \right) - \lambda \phi_1 \phi_2 \dots \phi_n$$

Free Field Dressing:

$$\Delta\mathcal{L}_2 = -\frac{\ell^2}{8} \left((\partial_\alpha \phi_i \partial^\alpha \phi_i)^2 - 2(\partial_\alpha \phi_i \partial^\alpha \phi_j)^2 + m_i^2 m_j^2 \phi_i^2 \phi_j^2 \right)$$



Reproduces $\mathcal{O}(\lambda\ell^2)$ -dressing
up to local polynomial terms

Hierarchy Problem

*Directly in terms of properties of the RG flow,
without ever mentioning quadratic divergencies*

For concreteness, let us place the discussion in the context
of non-SUSY GUTs

$$m_H \ll E \ll m_{GUT} : \mathcal{L} = CFT_{321} + \underset{\text{relevant}}{m_H^2 H^2} + \sum_i \frac{\mathcal{O}_i}{\underset{\text{irrelevant}}{M_{GUT}^{\Delta_i - 4}}}$$

How comes $m_H \ll m_{GUT}$ given no symmetry?

Hierarchy Problem

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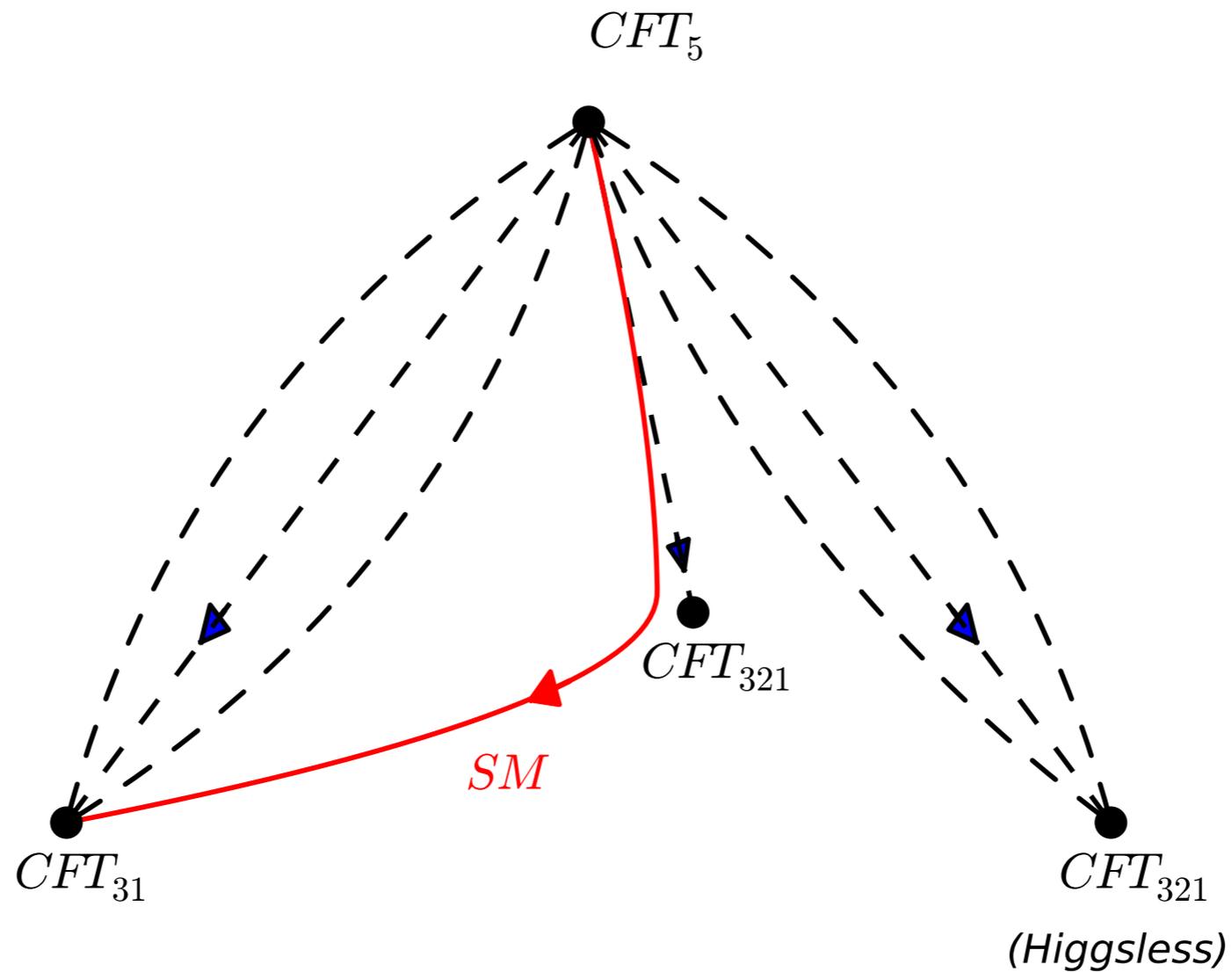
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How comes $m_H \ll m_{GUT}$ given no symmetry?

However, fine-tuning is truly manifest only as
seen from higher energies:

$$m_{GUT} \ll E : \mathcal{L} = CFT_5 + \underbrace{g_h m_{GUT}^2 H^2}_{\text{relevant}} + \underbrace{g_\Sigma m_{GUT}^2 \Sigma^2}_{\text{relevant}} + \dots$$



No picture like that in our example. Energy scale does not correspond to a threshold. No scale invariance and no Wilsonian RG above the scale.

Two notions of a naturalness:

- 1) *If a natural theory possesses unprotected relevant operators (scalar masses), the corresponding energy scale should be the highest*
- 2) *Among all possible scales set by **relevant** operators unprotected operators should correspond to the highest scale*

*Agree for QFT = UV CFT perturbed by relevant operators.

*May disagree in the presence of gravity.

Indeed, disagree in the gravitational dressing construction.

Is there a place for this scenario within the “standard” picture/string theory?



The moment we talk about naturalness we
are in the Landscape/Multiverse

Two canonical regions in the Landscape
capable of producing a light Higgs:

* An island where the Higgs mass is protected by a
symmetry (SUSY...)

* Among “ 10^{100} ” or so of random vacua with
randomly distributed values of the Higgs mass

Is there a third one?

* **Dragonland:** A (small) set of strongly coupled
vacua: $g_s = 1$ and Planckian extra dimensions

Possible lesson:

Should we be more serious about thinking on-shell when gravity is involved?

CC:

*Off-shell: nothing special about zero vacuum energy

*On-shell: zero CC is extremely special:

AdS:CFT, Minkowski:S-matrix, de Sitter: ???

Another possible lesson/alternative definition of naturalness:

Every natural QFT is an answer to some question.

Perhaps we should learn to ask more questions.

c.f. the following naturalness question:

31415926535897932384626433832795028841971693993...

is this sequence of digits “natural”?